

# Crystallization process of poly( $\epsilon$ -caprolactone)–poly(ethylene oxide)–poly( $\epsilon$ -caprolactone) investigated by infrared and two-dimensional infrared correlation spectroscopy

Jing Yu, Peiyi Wu\*

*Department of Macromolecular Science and the Key Laboratory of Molecular Engineering of Polymers, Fudan University, 220 Handan Road, 200433 Shanghai, PR China*

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## Abstract

The crystallization of poly( $\epsilon$ -caprolactone)–poly(ethylene oxide)–poly( $\epsilon$ -caprolactone) (PCL–PEO–PCL) triblock copolymer was studied using FTIR and 2D FTIR spectroscopies. The weight ratio of PCL/PEO in the investigated sample was about 20:1. Although it is such a low amount of PEO that it cannot form any crystals, the PEO block undergoes some structural change in the cooling process. It was established that the PCL constituent crystallized quickly, and then forced the noncrystallizable PEO constituent to form a tighter structure (helical conformation) from the *trans* zigzag conformation. Besides, through the 2D IR analysis, more exact and detailed assignments of the overlapped CH<sub>2</sub> bands have been made – the 1193 cm<sup>-1</sup> band is attributed to methylene next to the carbonyl group, whereas the 1162 cm<sup>-1</sup> and 1295 cm<sup>-1</sup> bands are assigned to other common methylenes.

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**Keywords:** Crystallization; FTIR; Two dimensional correlation analysis

## 1. Introduction

In recent years, a kind of amphiphilic triblock copolymer, poly( $\epsilon$ -caprolactone)–poly(ethylene oxide)–poly( $\epsilon$ -caprolactone) (PCL–PEO–PCL), has attracted much attention because of its environmental adaptability as well as potential capability as biomedical material. Both PCL and PEO constituents possess a number of outstanding physicochemical and biological properties which result in the fact that the copolymer mentioned above can be used as a drug delivery vehicle [1]. Furthermore, the hydrophilic PEO constituent is soluble in water while the hydrophobic PCL constituent can only dissolve in organic solvents. Thus, the copolymer PCL–PEO–PCL can form core–shell micelles through a self-assembly process in water.

A lot of research work has been made to investigate and promote the drug loading and delivery properties of the copolymer PCL–PEO–PCL. It was established that both biodegradability and drug permeability depend on the crystallinity [2]. However, there are only a few studies that have been published on the morphology and crystallization/melting properties of such block copolymer. Gan et al. [3] used WAXD, SAXS and DSC methods to investigate the isothermal crystallization and melting behavior of the PCL/PEO diblock copolymer. Besides the diblock copolymer, Bogdanov et al. [2] studied the isothermal crystallization behavior of three different types of PCL/PEG incompatible block copolymers including A–B, (A)<sub>2</sub>–B, A–B–A types by DSC. Floudas et al. [4] synthesized and studied the structure and dynamics of structure formation of a model triarm star block copolymer composed of PEO, PCL and PS by several methods including X-ray scattering, optical microscopy, AFM, DSC and rheology.

\* Corresponding author. Tel.: +86 21 65643255; fax: +86 21 65640293.  
E-mail address: [peiyiwu@fudan.edu.cn](mailto:peiyiwu@fudan.edu.cn) (P. Wu).

As mentioned above, a series of methods have been applied to study the crystallization behavior and properties of PCL/PEO block copolymers. But to the best of our knowledge, there are only few reports involving IR spectroscopy, a well-developed research method to investigate the conformation and local order of polymers. Since the real-time IR can give the information of structural evolution at the molecular level, the IR probe is widely used to reveal the structure and conformation changes during the crystallization and melting processes [5–8]. One of the reasons for the apparent lack of previous IR work may be due to the fact that the spectral region  $1300\text{--}1000\text{ cm}^{-1}$ , mainly associated with ester groups, is highly overlapped, which makes the detailed analysis and assignment rather difficult. However, the generalized two-dimensional (2D) correlation spectroscopy developed by Noda [9,10] in 1993 in addition to conventional IR methods may be able to solve this problem. The theory and introduction to 2D IR can be found elsewhere [9–12]. 2D IR introduced the FTIR to a second dimension, provided additional useful information and thus enhanced spectral resolution and pertinent features of spectral changes [13].

PEO and PCL are both crystallizable polymers, however, when they are combined into one copolymer as two blocks, the crystallization behavior only depends on the block length ratio. When the blocks length ratio is 3 or higher, only the longer block will crystallize [4]. We have investigated a series of PCL–PEO–PCL copolymer samples with different PEO/PCL weight ratios. In this work, we studied the triblock copolymer PCL–PEO–PCL with a PEO/PCL ratio of 1:20 and try to elucidate how the shorter noncrystallizable PEO block affects the crystallizing process of the PCL blocks. Although the PEO blocks in this sample may not exhibit a crystalline phase, they may still form some ordered conformations on a molecular scale. The helical conformation composed of a  $7/2$  helix and the *trans* zigzag conformation composed of a  $10/3$  helix are two of the most common structures that exist in the PEO polymer [14]. 2D IR is used to reveal the correlation between different groups, especially in the region of  $1300\text{--}1000\text{ cm}^{-1}$ . Samples with other ratios will be investigated in another paper, which is in preparation.

Quite different spectral patterns are found both in traditional IR and 2D IR analyses, which are believed to be produced by the dipole–dipole interaction between the C=O groups. Dipole–dipole interaction is a kind of inter- or intramolecular interaction between molecules or groups having a permanent electric dipole moment. The strength of the interaction depends on the strength and relative orientation of the dipoles, as well as on the distance between the centers and the orientation of the radius vector connecting the centers with respect to the dipole vectors [15]. Molecules will tend to align so that maximum attraction between dipoles is achieved. This occurs when the neighboring molecules have their dipoles anti-parallel, or the positive end of one dipole lies close to the negative end of another dipole.

## 2. Experimental details

### 2.1. Material and sample preparation procedures

The synthesis and purification of the triblock copolymer (PCL–PEO–PCL:  $M_w \approx 40,000\text{ g mol}^{-1}$ , PCL/PEO = 20:1 in weight) were performed according to procedures reported in the literature [1]. The thin copolymer films were cast on KBr windows from a 1% (w/v) PCL–PEO–PCL chloroform solution. The KBr windows were placed under a culture dish outside the fume cupboard to ensure the evaporating rate was slow enough, so that we could produce perfect films. The majority of the solvent evaporated after 24 h, and then the films were placed under vacuum at room temperature for 48 h to completely remove the residual solvent.

### 2.2. FTIR spectroscopy

For studying the isothermal crystallization process of the triblock copolymer PCL–PEO–PCL, the sample was set in a temperature controlling cell, which was placed in the sample compartment of the FTIR Spectrometer Nexus 470. The sample was first kept at  $80\text{ }^\circ\text{C}$  ( $T_m = 60\text{ }^\circ\text{C}$ ) until it was completely molten, then the temperature decreased at a rate of  $1.5\text{ }^\circ\text{C}/\text{min}$  to  $26\text{ }^\circ\text{C}$  and was kept at  $26\text{ }^\circ\text{C}$  for about 120 min. The FTIR measurements were carried out at  $26\text{ }^\circ\text{C}$ , in 2 min intervals. The spectra were collected at a resolution of  $4\text{ cm}^{-1}$  and 64 scans were accumulated. Before further analysis, all the spectra were baseline corrected.

### 2.3. Two-dimensional correlation analysis

The software for two dimensional correlation analysis in this present study is MATLAB7.0 (The Math Works Inc.). The correlation peaks are represented as pair  $(\nu_1, \nu_2)$  which refers, respectively, to the two independent wavenumbers involved. In all 2D correlation plots, the negative peaks are shaded while the remaining (unshaded) peaks are positive.

## 3. Results and discussion

### 3.1. Structure formation and bands assignment

The FTIR spectra of the investigated PCL–PEO–PCL sample at  $26\text{ }^\circ\text{C}$  (with comparatively high crystallinity) and at  $80\text{ }^\circ\text{C}$  (in the molten state) are shown in Fig. 1, reflecting large differences, especially in the region  $1500\text{--}800\text{ cm}^{-1}$ . The assignments of several absorption bands are summarized in Table 1. It should be noted that there are some small differences in the peak positions between literatures and our results. For example, according to Zhu et al. [16], the assignment of the C=O stretching region is very simple and clear. Nevertheless, since hydrogen bonding is not the dominating factor in our system, the peaks located at  $1724\text{ cm}^{-1}$  and  $1735\text{ cm}^{-1}$  are assigned, respectively, to the crystalline state and the amorphous regions. Fig. 2, in which the transformation can be seen clearly in the spectra obtained during the cooling process in

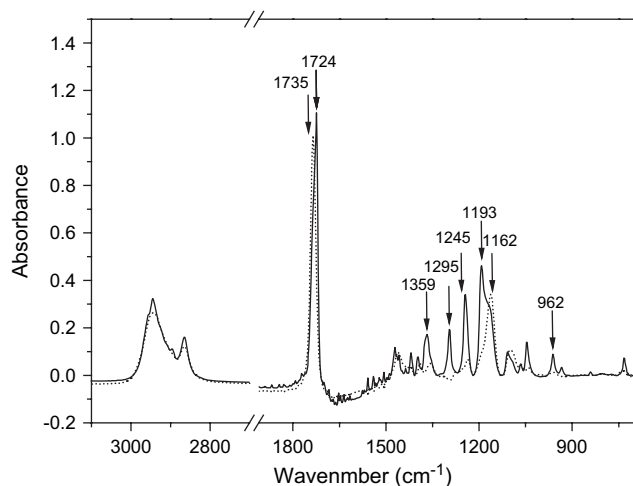


Fig. 1. FTIR spectra of the PCL-PEO-PCL sample in the molten (···) and crystalline states (—).

10 min interval, shows the change of the IR spectra around the C=O stretching region as a function of the cooling process. It is very obvious that during the cooling process, the band at  $1735\text{ cm}^{-1}$  decreases while the  $1724\text{ cm}^{-1}$  peak increases at the same time. Since the change of the C=O band can represent the conformational transformation of PCL constituent (crystallizable block of the investigated copolymer) to a large extent, we will consider it as a reference in the further 2D IR analysis. Moreover, as labeled in Fig. 2, there are many significant absorption bands sensitive to the conformation of the copolymer in the wavenumber range of  $1500\text{--}800\text{ cm}^{-1}$ . The peaks at  $1455\text{ cm}^{-1}$ ,  $962\text{ cm}^{-1}$ ,  $1295\text{ cm}^{-1}$  and  $1193\text{ cm}^{-1}$  which were attributed to  $\text{CH}_2$  scissoring [17],  $\text{CH}_2$  rocking [17],  $\text{CH}_2$  rocking combined with C–O–C deformation [17] in PEO block and two significant absorptions of crystalline PCL block, respectively [18], increase gradually as the temperature decreases, whereas absorption bands of another group at  $1359\text{ cm}^{-1}$  and  $1235\text{ cm}^{-1}$  increase accompanied by a slight shift to a higher frequency. However, there is only one peak at  $1162\text{ cm}^{-1}$  that shows significant attenuation during the

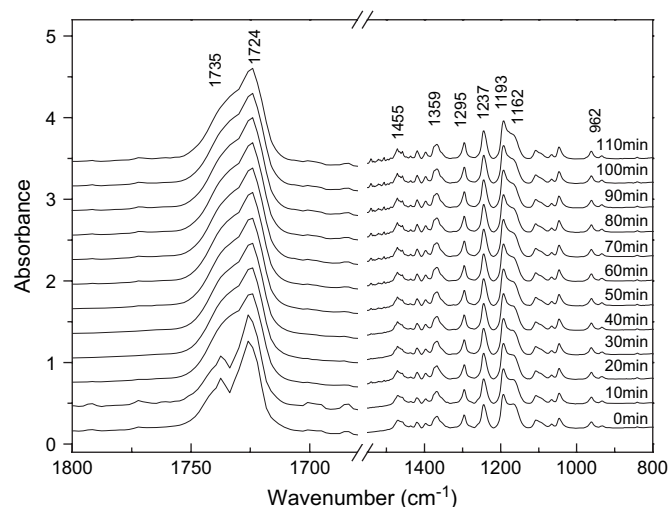


Fig. 2. Time-dependent FTIR spectra obtained at  $26\text{ }^\circ\text{C}$  of the investigated PCL-PEO-PCL sample in the region at around  $1700\text{ cm}^{-1}$  and  $1500\text{--}800\text{ cm}^{-1}$ .

cooling period, which is related to the amorphous phase in the PCL blocks according to literature [19].

The second-derivative IR spectra are also calculated for the amorphous state and the highly crystalline state. Several overlapped bands at around  $1245\text{ cm}^{-1}$ ,  $1455\text{ cm}^{-1}$ ,  $1067\text{ cm}^{-1}$  and  $1045\text{ cm}^{-1}$  can be found in both phases with a shift to a higher frequency in the crystalline state compared to the amorphous one.

To investigate the detailed mechanism of segmental motion taking place during the crystallization process, normalized peak heights of several characteristic peaks of different blocks are plotted as a function of cooling time, which is shown in Fig. 3. The  $1295\text{ cm}^{-1}$  and  $1193\text{ cm}^{-1}$  are both characteristic peaks of the PCL crystalline phase [18], the  $1367\text{ cm}^{-1}$  and  $1045\text{ cm}^{-1}$  absorptions are both considered as characteristic of a  $7/2$  helical conformation in the PEO blocks, while  $962\text{ cm}^{-1}$  and  $1245\text{ cm}^{-1}$  represent a  $10/3$  helical conformation [17]. It could be clearly seen in this picture that the crystallization process mainly took place within the period of 10 min after being cooled for an hour. Nevertheless, although the six bands

Table 1  
Band assignments of amorphous and crystalline PCL and PEO blocks [16–19]

Band	Wavenumber ( $\text{cm}^{-1}$ )		Assignment
	PCL	PEO	
C=O	1736 [16]		C=O stretching (amorphous phase)
	1724 [16]		C=O stretching (crystalline phase)
C–H		1468 [17]	$\text{CH}_2$ scissoring
		1359 [17]	$\text{CH}_2$ wagging
		1242 [17]	$\text{CH}_2$ twisting and C–O–C stretching
		1060 [17]	$\text{CH}_2$ rocking and C–O–C stretching
		964 [17]	$\text{CH}_2$ rocking
		1162 [19]	C–O–C stretching of PCL amorphous phase
C–O–C		1242 [17]	$\text{CH}_2$ twisting and C–O–C stretching
		1060 [17]	$\text{CH}_2$ rocking and C–O–C stretching
	1295 [18]		Significant absorption of PCL crystalline phase
	1192 [18]		Significant absorption of PCL crystalline phase

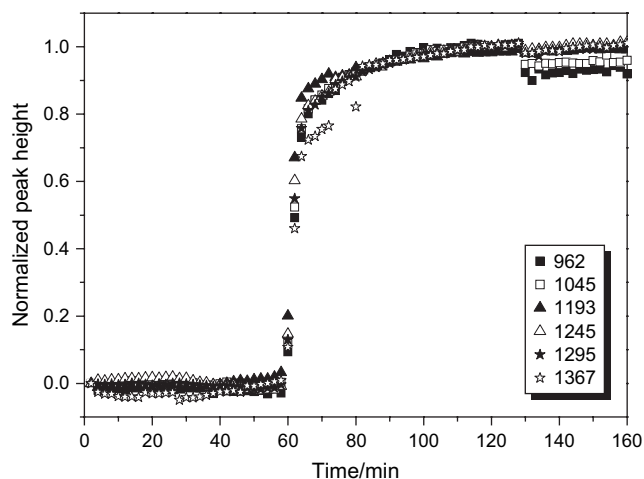


Fig. 3. Normalized peak heights of the characteristic peaks of the PCL crystalline phase at 1295  $\text{cm}^{-1}$  and 1193  $\text{cm}^{-1}$  and of different PEO conformation at 1367  $\text{cm}^{-1}$ , 1245  $\text{cm}^{-1}$ , 1045  $\text{cm}^{-1}$  and 962  $\text{cm}^{-1}$  as a function of time of the isothermal crystallization process calculated from the spectra shown in Fig. 2.

are attributed, respectively to different conformations of the PCL or PEO blocks, there are hardly any differences between the lines in Fig. 3. Obviously, the one-dimensional analysis is not able to provide sufficient evidence. We use in this study 2D IR correlation spectroscopy to elucidate the detailed crystallization mechanism of the block copolymer.

### 3.2. Two-dimensional FTIR analysis

The synchronous and asynchronous correlation spectra of PCL–PEO–PCL of the C=O stretching region are shown in Fig. 4. The two autopeaks shown in the synchronous spectra locate at (1736  $\text{cm}^{-1}$ , 1736  $\text{cm}^{-1}$ ) and (1724  $\text{cm}^{-1}$ , 1724  $\text{cm}^{-1}$ ), respectively. Since the 2D calculations always obtain average results, the center of the cross peak may slightly shift from the

peak in 1D reference spectra. The negative area in Fig. 4A centers at (1740  $\text{cm}^{-1}$ , 1724  $\text{cm}^{-1}$ ). Thus, there are obviously two kinds of C=O bands existing in the IR spectrum, one of which can be found in the region of a higher wavenumber at around 1740  $\text{cm}^{-1}$  while the other at a lower wavenumber at around 1724  $\text{cm}^{-1}$ . Furthermore, the 2D analysis also indicates that the peak appearing at higher wavenumber alters in the opposite direction to the peak at lower one, similar to the observation in the 1D IR analysis.

The negative asynchronous cross peak  $\Psi(1735 \text{ cm}^{-1}, 1724 \text{ cm}^{-1})$  developed in Fig. 4B indicates that the intensities of two dynamic bands at 1735  $\text{cm}^{-1}$  and 1724  $\text{cm}^{-1}$  vary out of phase, and provides the information on the sequential order of intensity changes between these two peaks. According to the publications of Noda [9], if the product of the intensity of a synchronous peak and its corresponding asynchronous peak is positive (a concise expression of this sentence is:  $\Phi(\nu_1, \nu_2) \cdot \Psi(\nu_1, \nu_2) > 0$ ), band  $\nu_1$  varies prior to band  $\nu_2$ . Contrarily, if  $\Phi(\nu_1, \nu_2) \cdot \Psi(\nu_1, \nu_2) < 0$ , band  $\nu_2$  varies prior to the band  $\nu_1$ . It should be mentioned that all peaks in this paper are the peaks appearing in the top-left region of the maps, except for those maps without a top-right to left-bottom symmetry axis. The cross peaks  $\Phi(1735 \text{ cm}^{-1}, 1724 \text{ cm}^{-1})$  and  $\Psi(1735 \text{ cm}^{-1}, 1724 \text{ cm}^{-1})$  are both negative, that is, the product of these two peaks are positive ( $\Phi(1735 \text{ cm}^{-1}, 1724 \text{ cm}^{-1}) \cdot \Psi(1735 \text{ cm}^{-1}, 1724 \text{ cm}^{-1}) > 0$ ), therefore the band at 1735  $\text{cm}^{-1}$  varies prior to the band at 1724  $\text{cm}^{-1}$ . Moreover, similar to Zhang et al.'s study [5], there is a so-called “butterfly pattern” in the C=O stretching region in the asynchronous correlation spectra. Usually, the appearance of this pattern is attributed to a peak shift combined with the intensity changes or band splitting [20,21].

Fig. 5 shows the 2D correlation spectra of the highly overlapped 1380–900  $\text{cm}^{-1}$  region. In principle, we should be able to extract useful information about the crystallizing process of the copolymer system from these data. Among many cross peaks shown in Fig. 5A (synchronous spectrum), only the

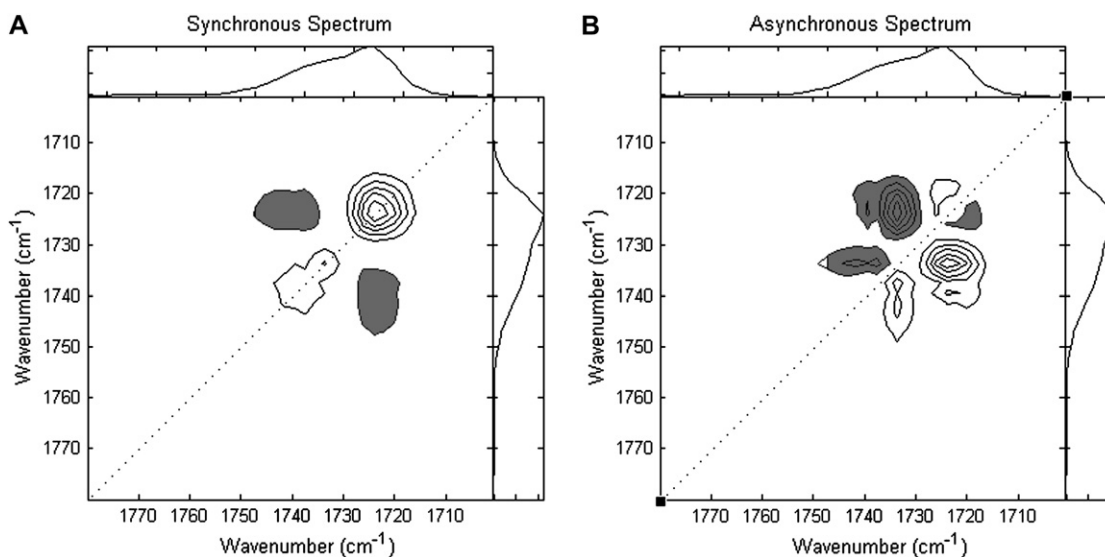


Fig. 4. Synchronous and asynchronous 2D correlation spectra of PCL–PEO–PCL copolymer at 1780–1700  $\text{cm}^{-1}$ .

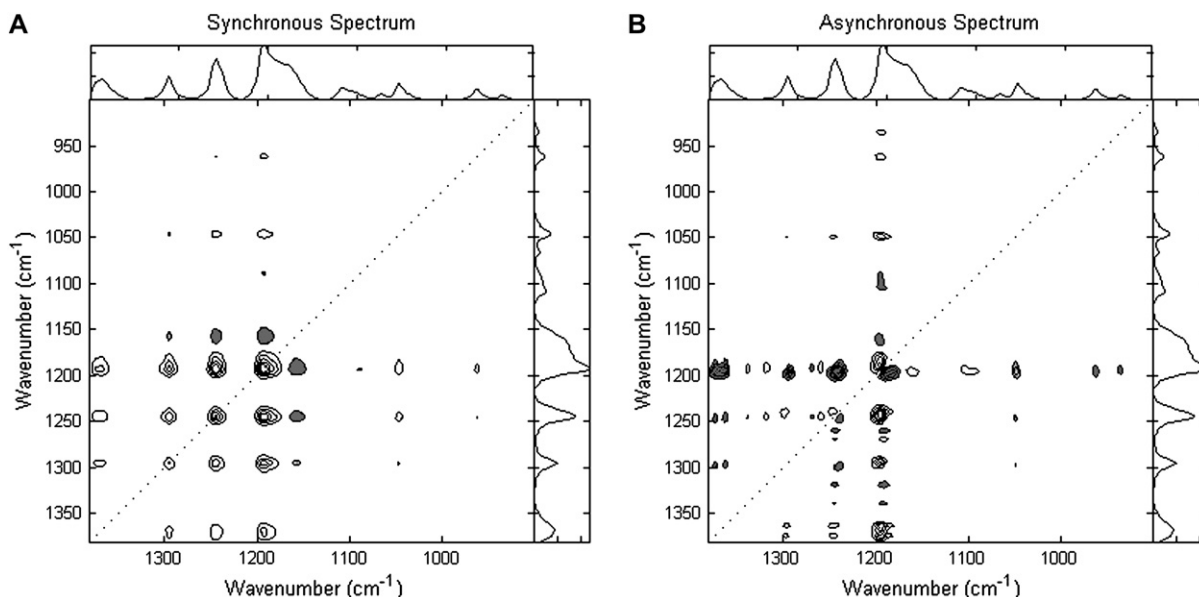


Fig. 5. Synchronous and asynchronous 2D correlation spectra of PCL-PEO-PCL copolymer at 1380–900  $\text{cm}^{-1}$ .

cross peaks at 1162  $\text{cm}^{-1}$  are negative, which suggests that the majority of peaks in the range of 1380–900  $\text{cm}^{-1}$  alter in the same direction, whereas the peak at 1162  $\text{cm}^{-1}$  changes oppositely. As outlined before, the 1162  $\text{cm}^{-1}$  band is assigned to the amorphous PCL phase [19]. Based on the above analysis, it is easy to understand that the peaks altering in the opposite direction to the band at 1162  $\text{cm}^{-1}$  should be attributed to the formation of a more regular conformation. However, the PEO block cannot take the conformation of the crystalline phase because the block length ratio of PCL/PEO is much higher than 3 [4].

It is, however, somewhat difficult to pick out the fine structure of the spectrum in Fig. 5, when such a broad spectral region is displayed. The full advantage of the contour map

representation of 2D correlation spectra will not be realized until the technique is performed on a limited portion of the spectral map. Parts A and B of Fig. 6 are the close-up view of the contour maps shown in Fig. 5 in the spectral region between 1380  $\text{cm}^{-1}$  and 1230  $\text{cm}^{-1}$  in one dimension and between 1280  $\text{cm}^{-1}$  and 1170  $\text{cm}^{-1}$  in another dimension. Thus, more detailed features become apparent.

After the analysis of the asynchronous correlation spectrum in Fig. 6B, several cross peaks (denoted as a, b and c) are found to contain interesting information. The positions of these three peaks are as follows: a (1367  $\text{cm}^{-1}$ , 1245  $\text{cm}^{-1}$ ), b (1295  $\text{cm}^{-1}$ , 1193  $\text{cm}^{-1}$ ) and c (1295  $\text{cm}^{-1}$ , 1245  $\text{cm}^{-1}$ ). Peaks (a) and (b) are negative, whereas (c) is positive. The corresponding synchronous cross peaks are all positive. Following

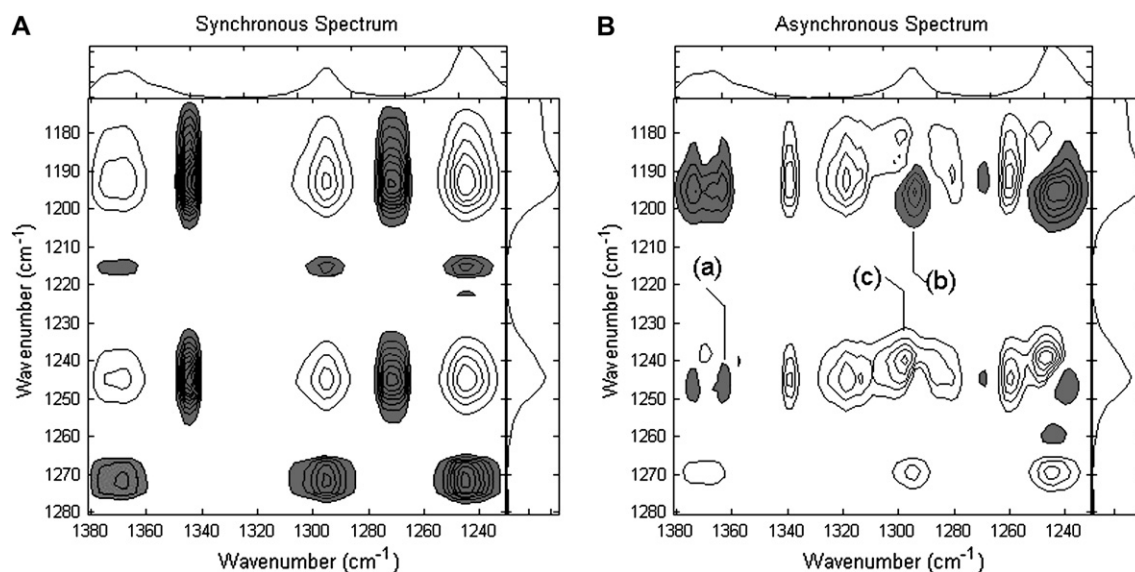


Fig. 6. The signs of cross peaks appearing in the expanded view of the contour maps shown in Fig. 5.

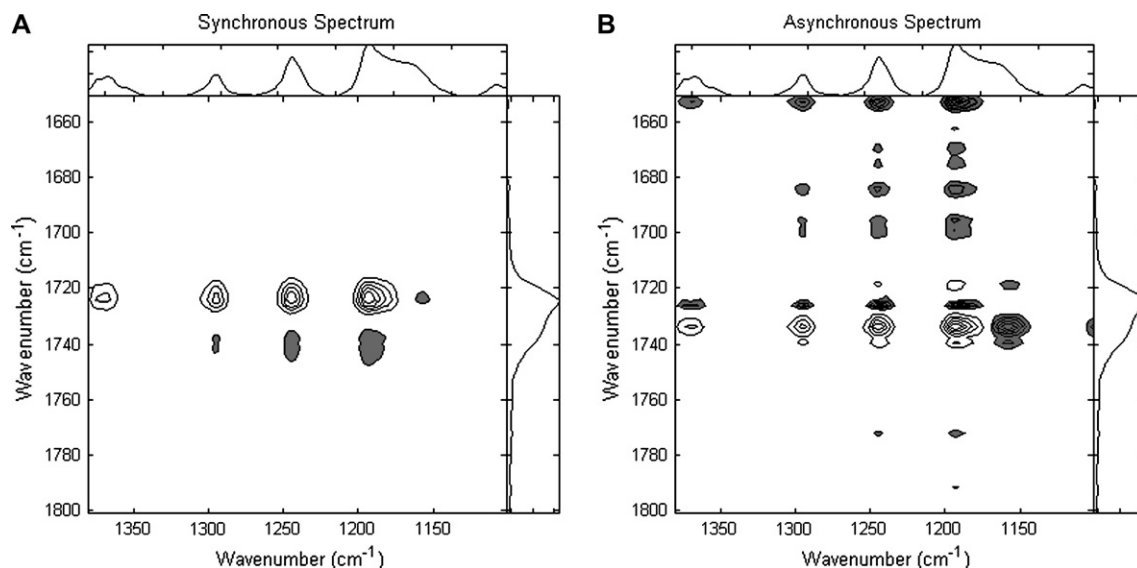


Fig. 7. Synchronous and asynchronous 2D FTIR correlation spectra of PCL–PEO–PCL copolymer obtained during the isothermal crystallization process between 1380–1100  $\text{cm}^{-1}$  and 1800–1650  $\text{cm}^{-1}$ .

the rule mentioned above, a conformational transformation sequence among the bands involved is obtained: 1193  $\text{cm}^{-1}$  > 1295  $\text{cm}^{-1}$  > 1245  $\text{cm}^{-1}$  > 1367  $\text{cm}^{-1}$ . The significance of this sequence notation will be discussed later.

Additionally, 2D correlation spectra between different ranges are acquired to achieve information about the relationship between different groups in the copolymer. The 2D spectra between the C=O stretching region and the highly overlapped region mainly associated with ester groups are shown in Fig. 7. In the synchronous spectra (Fig. 7A), all peaks involving the 1724  $\text{cm}^{-1}$  band are positive and those involving the 1735  $\text{cm}^{-1}$  band are negative, except for the

1162  $\text{cm}^{-1}$  band. From the asynchronous spectrum, we inferred that a band split occurs in the C=O stretching mode, which always represent a dipole–dipole coupling when the molecules form an ordered structure [5]. This result goes consistent with what we get from the self-correlation analysis of the C=O stretching region. Also, the changing sequence can be obtained easily so that the bands at 1735  $\text{cm}^{-1}$  and 1724  $\text{cm}^{-1}$  vary prior to all the peaks in the range of 1380–1100  $\text{cm}^{-1}$ .

To investigate the role of different C–H groups in the crystallization process, 2D analysis between regions containing information of C–H and C–O–C bands is performed. Fig. 8

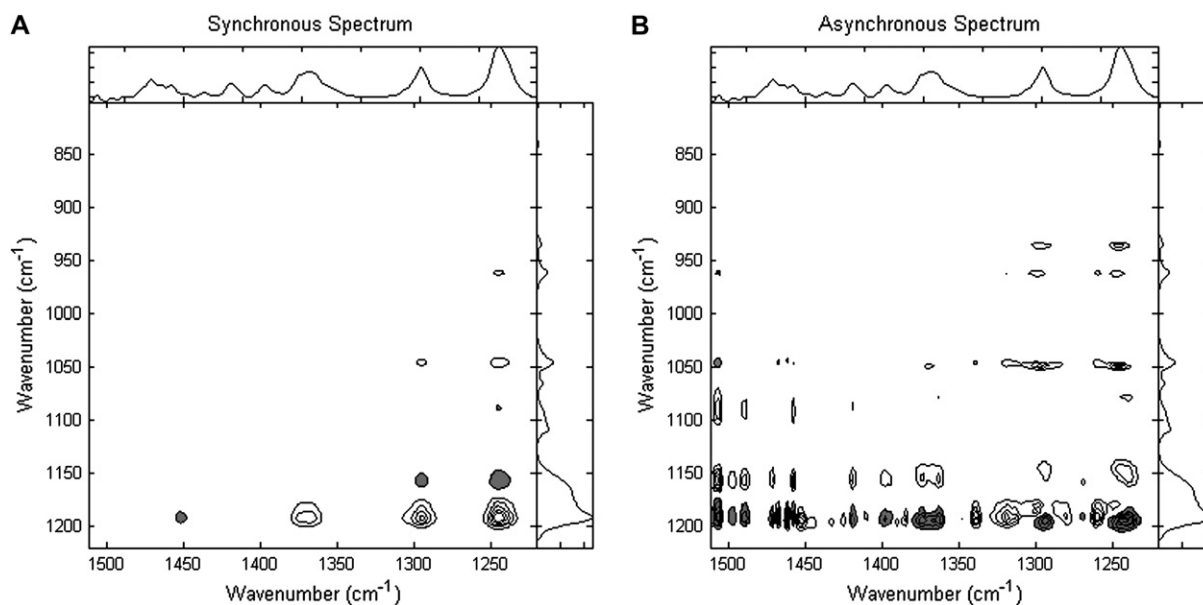


Fig. 8. Synchronous and asynchronous 2D FTIR correlation spectra of the investigated PCL–PEO–PCL copolymer obtained during the isothermal crystallization process between 1210–800  $\text{cm}^{-1}$  and 1510–1220  $\text{cm}^{-1}$ .

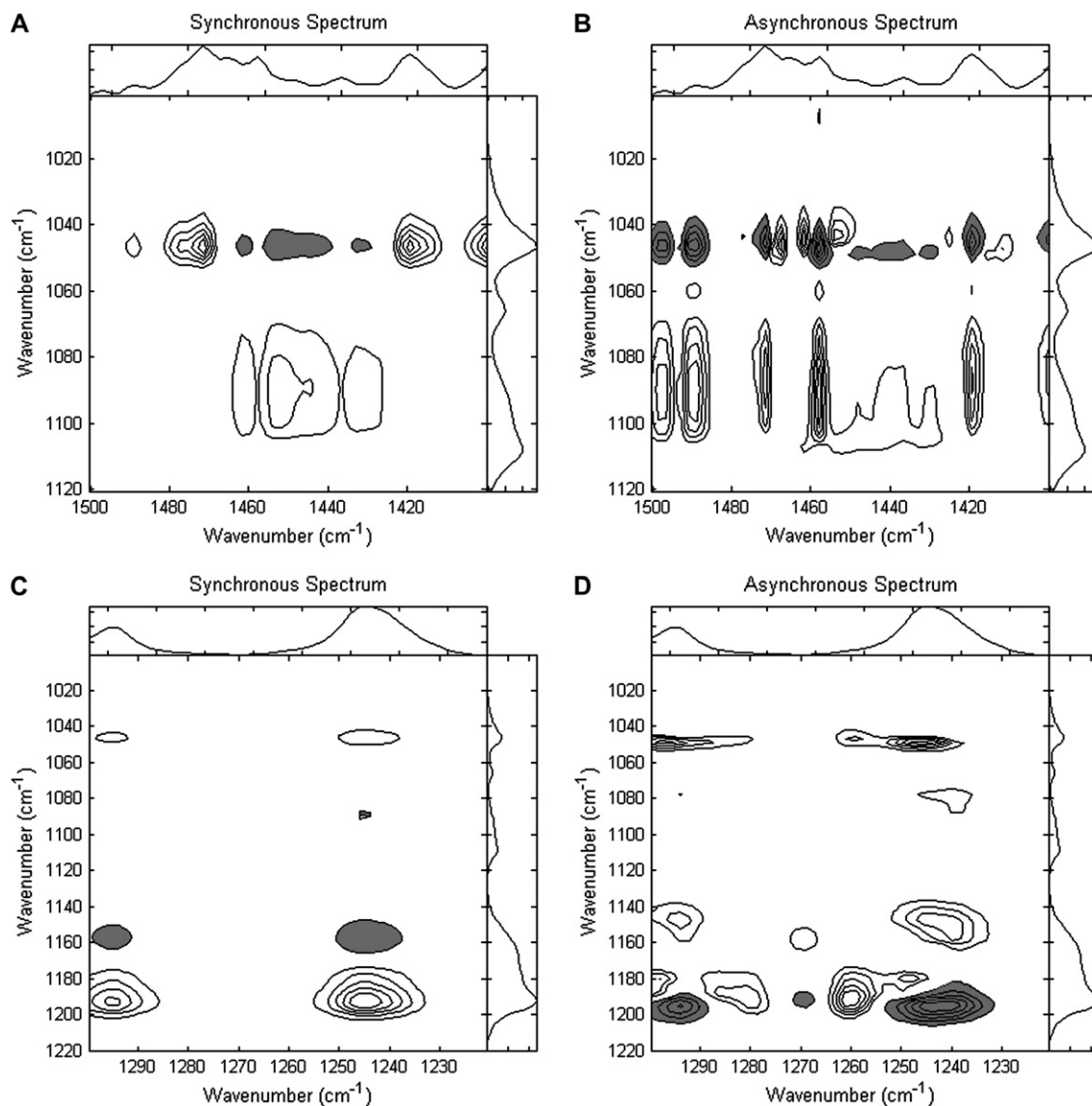
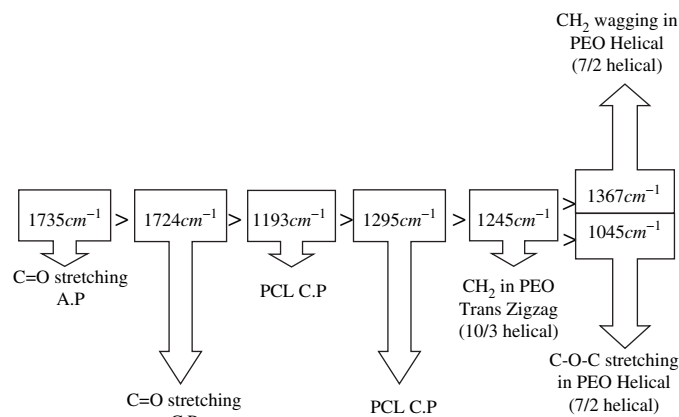


Fig. 9. The signs of cross peaks appearing in the expanded view of the contour maps shown in Fig. 8.

shows the 2D correlation spectra between 1210–800  $\text{cm}^{-1}$  and 1510–1220  $\text{cm}^{-1}$ . Since the whole region is too large for the 2D calculation to get a satisfactory result, several localized views are made in addition as shown in Fig. 9. The parts A and B in Fig. 9 are synchronous and asynchronous 2D IR spectra between 1120–1000  $\text{cm}^{-1}$  and 1500–1400  $\text{cm}^{-1}$ , while the parts C and D represent the regions between 1220–1020  $\text{cm}^{-1}$  and 1300–1220  $\text{cm}^{-1}$ . Three meaningful cross peaks can be found in the two asynchronous spectra (Fig. 9B and D) at the position of (1452  $\text{cm}^{-1}$ , 1045  $\text{cm}^{-1}$ ), (1295  $\text{cm}^{-1}$ , 1045  $\text{cm}^{-1}$ ) and (1245  $\text{cm}^{-1}$ , 1045  $\text{cm}^{-1}$ ). The first one is negative while the other two are positive. And the corresponding synchronous cross peaks are negative, positive and positive, respectively. Thus, another band variation sequence can be achieved: 1045  $\text{cm}^{-1}$  varies after the other three bands at 1295  $\text{cm}^{-1}$ , 1452  $\text{cm}^{-1}$  and 1245  $\text{cm}^{-1}$ .

Summing up the conformational transformation sequences obtained during the 2D correlation analysis, a final

result is attained which can be described by the following scheme:



Here, A.P. means amorphous phase whereas C.P. means crystalline phase.

In this sequence,  $1735\text{ cm}^{-1}$  and  $1724\text{ cm}^{-1}$  are attributed to the C=O stretching band in the amorphous and crystalline phases, respectively [16]. The bands at  $1193\text{ cm}^{-1}$  and  $1295\text{ cm}^{-1}$  are related to absorptions of the PCL crystalline phase [18]. The 2D results revealed the information that we cannot get from the 1D analysis shown in Fig. 3. The band at  $1245\text{ cm}^{-1}$  is assigned to a CH<sub>2</sub> twisting vibration of PEO in the *trans* zigzag conformation [17]. The band at  $1367\text{ cm}^{-1}$  is attributed to a CH<sub>2</sub> wagging motion in the PEO helical conformation [17]. The band at  $1060\text{ cm}^{-1}$ , which appears at  $1045\text{ cm}^{-1}$  in our polymer system, is also attributed to the helical conformation but mainly relates to the C–O–C stretching vibration [17].

According to the above data, we can derive on a molecular level that, in the cooling process, the structural adjustment of the PCL blocks occurs faster than that of the PEO blocks, which means that the PCL blocks are able to form a crystalline phase more easily. The detailed mechanism in each one of these two kinds of blocks are: in PCL blocks, the carbonyl groups alter faster than the CH<sub>2</sub> groups, and there are also some differences between different CH<sub>2</sub> groups in the PCL blocks; in the PEO blocks, the CH<sub>2</sub> groups vary prior to the C–O–C groups, and what is more important, a conformational transformation is uncovered which indicates that the PEO segments change from an incompact conformation into a tighter one.

We consider that it is the dipole–dipole interaction between the carbonyl groups that cause the C=O groups change prior to others. To prove this interaction actually exists in our system and is strong enough to affect the crystallization process of the block copolymer, we have two evidences. Firstly, as has been mentioned in several literatures references [6,22], the dipole–dipole interactions always result in a band splitting for polymer systems when the molecules form an ordered structure. Furthermore, the high-frequency shift of the C=O stretching band is usually due to the dipole coupling of the C=O bonds in the crystalline state. From both the traditional spectra (Fig. 2) and the 2D correlation analysis of C=O stretching region, we can derive that in our system, there is obviously a band splitting taking place together with a high-frequency shift of the C=O stretching band. Secondly, the carbonyl group has the largest dipole moment of almost all chemical bonds. The value of the dipole moment of the carbonyl group is 2.3–2.7, whereas the one of the O–H bond in water is 1.85, which determines that the dipole–dipole interaction between carbonyl groups is stronger than the hydrogen-bonding effect. These two evidences substantiate the existence and the possible effect of dipole–dipole interaction in our system, which is very helpful to explain the mechanism brought forward through our analysis. The carbonyl groups with strong polarity have a comparatively strong trend to congregate first and produce the driving force for the neighboring molecular chain to form an ordered state. It is obvious that, in this process, the CH<sub>2</sub> groups located near the carbonyl groups will be affected more easily compared to those which are located further away from the carbonyl groups. With reference to the transformational sequence we obtained from the two-dimensional analysis of both the bands at  $1193\text{ cm}^{-1}$  and

$1295\text{ cm}^{-1}$ , which were both attributed to absorptions of the crystalline PCL phase, we suppose that the  $1193\text{ cm}^{-1}$  band can be attributed to the methylene group next to the carbonyl group, whereas the  $1295\text{ cm}^{-1}$  band should be assigned to other common methylene groups.

As we have mentioned in the analysis of Fig. 3, the crystallization process occurs very quickly, which indicates that the crystallizable segments of the PCL block form regular structure in a very short period of time. In this process, the noncrystallizable segments of the PEO block are separated between the ordered PCL segments, and will be certainly influenced by the PCL crystallization process. It is probably this effect that influences the conformational transition of the PEO blocks.

#### 4. Conclusion

FTIR and 2D FTIR spectroscopies have been applied to study the crystallization behavior of a PCL-rich PCL–PEO–PCL copolymer. The PEO blocks are in such a low amount that they cannot perform crystallization. However, we investigated the role of the PEO blocks in the copolymer and put forward hypothesis of the crystallization process. The carbonyl groups of the PCL blocks first congregate and force the neighboring chains to form an ordered state. The squeezing effect produced during the crystallization process of the PCL blocks forces the PEO blocks to form a tighter and more ordered state, which appears as the transformation from the *trans* zigzag (10/3 helical) to the helical conformation (7/2 helical). Besides, some more exact and detailed assignments have been made. The  $1193\text{ cm}^{-1}$  band is attributed to methylene next to the carbonyl group, whereas  $1295\text{ cm}^{-1}$  bands are assigned to other common methylenes.

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